## BRIEF COMMUNICATIONS

THE ANALYTIC REPRESENTATION OF THE PULSE

## ENERGY IN THE WAKE

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From the equation of balance for the pulse energy we find the distribution of the average turbulent energy in the automodel region of a two-dimensional wake.

The equations of balance for the momentum and the kinetic energy of turbulence in a remote part of a two-dimensional wake in an incompressible fluid can be written as

$$
\begin{gather*}
\left.U_{1} \frac{\partial U}{\partial x_{1}}=-\frac{\partial}{\partial x_{2}} \overline{\left(v_{2}^{\prime} v_{1}^{\prime}\right.}\right) \\
U_{1} \frac{\partial E}{\partial x_{1}}=-\frac{\partial}{\partial x_{2}}\left[\overline{\left.v_{2}^{\prime} \frac{1}{2} \sum_{i=1}^{3} v_{i}^{\prime 2}+\frac{1}{\rho} \overline{v_{2}^{\prime} p^{\prime}}\right]-\overline{v_{2}^{\prime} v_{1}^{\prime}} \frac{\partial U}{\partial x_{2}}-v \sum_{i, \alpha=1}^{3} \overline{\left(\frac{\partial v_{i}^{\prime}}{\partial x_{\alpha}}\right)^{2}},}\right. \tag{1}
\end{gather*}
$$

where $E=\sum_{i=1}^{3} \overline{v_{i}^{!2}} / 2$; the bar denotes the time average.
We can express the turbulent friction in Prandtl's form:

$$
\begin{equation*}
\overline{v_{2}^{\prime} v_{1}^{\prime}}=-v_{\mathrm{T}} \frac{\partial U}{\partial x_{2}}, v_{\mathrm{T}}=\left(U_{1}-U_{0}\right) \delta \tag{2}
\end{equation*}
$$

and put, for example, as in [1]

$$
\begin{equation*}
v_{2}^{\prime} \frac{1}{2} \sum_{i=1}^{3} v_{i}^{\prime 2}+\frac{1}{\rho} \overline{v_{2}^{\prime} p^{\prime}}=-v_{\mathrm{T}} \frac{\partial E}{\partial x_{2}}, v \sum_{i, \alpha=1}^{3} \overline{\left(\frac{\partial v_{i}^{\prime}}{\partial x_{\alpha}}\right)^{2}}=c v_{\mathrm{T}} \frac{E}{\delta^{2}} \tag{3}
\end{equation*}
$$

where c is a constant to be determined experimentally. Introducing the new variables

$$
\begin{aligned}
u & =U_{1}-U, \frac{u}{u_{0}}=f(\eta), \eta=\frac{x_{2}}{\delta}, \\
\frac{u_{0}}{U_{1}} & =\psi\left(x_{1}\right), \frac{E}{E_{0}}=h(\eta), \quad \frac{E_{0}}{U_{1}^{2}}=\varphi\left(x_{1}\right)
\end{aligned}
$$

and using the integral condition in the form

$$
\begin{equation*}
\psi \delta=\frac{c_{x} d}{4} \frac{1}{\int_{0}^{\infty} f d \eta}=\text { const }, \tag{4}
\end{equation*}
$$

we find that there is similarity if

$$
\begin{gather*}
\delta^{\prime} / \psi=-\psi^{\prime} \delta / \psi^{2}=-\varphi^{\prime} \delta / 2 \varphi \psi=B=\text { const }  \tag{5}\\
\psi^{2} / \varphi=D=\text { const. }
\end{gather*}
$$

Then (1) takes the form

$$
\begin{equation*}
f^{\prime \prime}+z f^{\prime}+f=0 \tag{6}
\end{equation*}
$$

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Fig. 1. Experimental data from [2] (continuous curve) and computation (dotted curve).

$$
\begin{equation*}
h^{\prime \prime}+z h^{\prime}+(2-c / B) h=-D f^{\prime 2} \tag{7}
\end{equation*}
$$

where the primes denote differentiation with respect to $z=\sqrt{B} \eta$.
We know the solution of Eq. (6) with the following boundary conditions:

$$
f=1 \text { for } z=0, f \rightarrow 0 \text { as } z \rightarrow \infty
$$

to be

$$
\begin{equation*}
f=\exp \left(-z^{2} / 2\right) \tag{8}
\end{equation*}
$$

The pulse-energy profile must satisfy the boundary conditions

$$
\begin{equation*}
h=1 \text { for } z=0, h \rightarrow 0 \text { as } z \rightarrow \infty, \tag{9}
\end{equation*}
$$

and also the symmetry condition

$$
\begin{equation*}
h^{\prime}=0 \text { for } z=0 \tag{10}
\end{equation*}
$$

We can show that we can describe the experimental results completely satisfactorily by taking $c=3 B$. Then the solution of Eq. (7), taking account of (8), with the boundary conditions (9) can be written as

$$
\begin{equation*}
h=z\{(3 / 2) D \sqrt{\pi}[1-\Phi(z \sqrt{2})]-(1+2 D) \sqrt{\pi / 2}[1-\Phi(z)]\}+(1+2 D) \exp \left(-z^{2} / 2\right)-2 D \exp \left(-z^{2}\right) \tag{I1}
\end{equation*}
$$

where

$$
\Phi(x) \equiv \sqrt{2 / \pi} \int_{0}^{x} \exp \left(-t^{2} / 2\right) d t
$$

It follows from this that

$$
h^{\prime}(0)=(1+2 D) \sqrt{\pi / 2}-3 D_{\sqrt{2}} \sqrt{\pi} / 2 .
$$

Using (10), we have

$$
D=\frac{\sqrt{2}}{3-2 \sqrt{2}} \approx 8.2
$$

In Fig. 1 there is a comparison between the computed and the experimental values [2]. The ordinate is, from (5), (11):

$$
\begin{equation*}
\frac{2 E}{u_{0}^{2}}=\frac{2}{D} \dot{h}=3 \sqrt{\pi} z[\Phi(z)-\Phi(z \sqrt{2})]+(2 / D+4) \exp \left(-z^{2} / 2\right)-4 \exp \left(-z^{2}\right) \tag{12}
\end{equation*}
$$

The abscissa is $\xi=x_{2} / \sqrt{d} x_{1}$, and we can find that $z=5.42 \xi$, since, by (4) and (5),

$$
\delta / \sqrt{B}=\sqrt{c_{x} \sqrt{B} d x_{1} / 2}
$$

and, by experiment [2], $\sqrt{\mathrm{B}}=1 / \mathrm{R}_{\mathrm{T}}=0.08, \mathrm{c}_{\mathrm{X}}=0.85$.
Using (8) and (12) it is easy to compute the ratio of the total intensities of the change in the average velocity and pulse velocity in the automodel region:

$$
\frac{\int_{-\infty}^{\infty} u^{2} d x_{2}}{\int_{-\infty}^{\infty} 2 E d x_{2}}=\frac{D}{2} \frac{\int_{0}^{\infty} f^{2} d \eta}{\int_{0}^{\infty} h d \eta}=2
$$

which agrees exactly with experiment (cf. [3], p. 173, Fig. 7, 1b).

## NOTATION

| $U$ | is the longitudinal component of average velocity; |
| :--- | :--- |
| $\mathrm{U}_{1}$ | is the velocity of the undisturbed flow; |
| $\mathrm{v}_{\mathrm{i}}^{\prime}, \mathrm{p}^{\prime}$ | are the velocity and pressure pulsations; |

$y$ is the kinematic viscosity coefficient;
$\delta \quad$ is the relative wake width;
$\rho \quad$ is the density;
d is the height of middle section of body;
$c_{\mathrm{X}} \quad$ is the drag coefficient.

## Subscript

0 denotes the wave axis.

## LITERATURE CITED

1. K. E. Dzhaugashtin, Magnitnaya Gidrodinamika, 4, 64 (1968).
2. O. Khintse, Turbulence [in Russian], Fizmatgiz, $\bar{M}$ Moscow (1963).
3. A. A. Townsend, The Structure of Turbulent Flow with Transverse Shear [Russian translation], IL (1959).
